

THRESHOLD IRRADIATION DOSE FOR AMORPHIZATION OF SILICON CARBIDE — L. L. Snead and S. J. Zinkle (Oak Ridge National Laboratory)

OBJECTIVE

The amorphization of silicon carbide under neutron irradiation is a concern at low irradiation temperatures. The objective of this work is to present a summary of the data on amorphization due to ion beam irradiation and to propose a model which describes the damage threshold as a function of temperature for ion and neutron irradiated material.

SUMMARY

The amorphization of silicon carbide due to ion and electron irradiation is reviewed with emphasis on the temperature-dependent critical dose for amorphization. The effect of ion mass and energy on the threshold dose for amorphization is summarized, showing only a weak dependence near room temperature. Results are presented for 0.56 MeV silicon ions implanted into single crystal 6H-SiC as a function of temperature and ion dose. From this, the critical dose for amorphization is found as a function of temperature at depths well separated from the implanted ion region. Results are compared with published data generated using electrons and xenon ions as the irradiating species. High resolution TEM analysis is presented for the Si ion series showing the evolution of elongated amorphous islands oriented such that their major axis is parallel to the free surface. This suggests that surface or strain effects may be influencing the apparent amorphization threshold. Finally, a model for the temperature threshold for amorphization is described using the Si ion irradiation flux and the fitted interstitial migration energy which was found to be ~ 0.56 eV. This model successfully explains the difference in the temperature-dependent amorphization behavior of SiC irradiated with 0.56 MeV silicon ions at 1×10^{-3} dpa/s and with fission neutrons irradiated at 1×10^{-6} dpa/s irradiated to 15 dpa in the temperature range of $\sim 340 \pm 10$ K.

PROGRESS AND STATUS

Introduction

An understanding of the microstructural evolution of SiC under irradiation is critical to the application of these materials in fusion systems. At the expected application temperatures for SiC ($< 1200^\circ\text{C}$) the most significant microstructural changes occur in the lower temperature regime. Specifically, at temperatures under 100°C , volumetric swelling due to point defect induced strain has been seen to reach 3%.^{1,2} At these low temperatures, amorphization of the SiC is also possible, which would lead to a substantial volumetric expansion of $\sim 15\%$,³⁻⁵ along with correspondingly large changes in mechanical properties such as hardness and modulus.^{3,4,6} Specifically, Snead et al.⁶ found that the elastic modulus for chemically vapor infiltrated SiC decreased from 400 GPa in the unimplanted region, to ~ 275 GPa in the amorphized region. Likewise, the hardness was seen to decrease from 41 GPa to 32 GPa. Similar behavior has been recently reported by other authors⁵ for ion irradiated SiC.

Numerous studies^{3, 7-14} have shown that SiC becomes amorphous during ion-beam irradiation at temperatures between 77 K and room temperature for damage levels of approximately 0.1 to 0.5 atomic displacements per atom (dpa.) To date there has been no demonstration of neutron induced amorphization of SiC, although it is worth noting that there is a lack of low temperature high dose neutron irradiation data. The majority of studies relating to SiC amorphization have utilized low-energy (< 1 MeV) ion-beams⁸⁻¹⁴ with a few researchers studying the amorphization threshold using high energy electrons.¹⁵⁻¹⁹ These electron and ion irradiation studies have typically been carried out at temperatures between liquid nitrogen and 470 K. Amorphization was not observed in SiC irradiated with fission neutrons (damage rate $\sim 10^{-7}$ dpa/s) at ~ 420 K to damage levels in excess of 1 dpa.²⁰

In the past ten years there has been a growing interest in the effect of irradiation temperature on the amorphization of silicon carbide. Specifically, Kinoshita¹⁵ and Matsunaga¹⁶ used 1 MeV electron irradiation to measure the damage threshold for amorphization from liquid nitrogen temperature to room temperature. Inui¹⁷⁻¹⁹ has used 2 MeV TEM electrons and Weber^{5, 21-24} has used a 1.5 MeV Xe beam in conjunction with (non damaging) TEM to study amorphization. In all cases the critical threshold for amorphization was measured as a function of temperature. It was observed that there is a temperature-independent amorphization dose at low temperatures followed by a temperature above which the damage level required to amorphize SiC increases rapidly. In each case an apparent asymptotic increase in the critical dose for amorphization occurred, yielding a critical temperature above which amorphization is apparently not possible. This value is reported to be significantly lower for electron irradiations compared to ion irradiations at the same damage rate of $\sim 1 \times 10^{-3}$ dpa/s. The critical temperatures above which amorphization does not occur has been reported to range between 290 K and 340 K for electrons,¹⁷⁻¹⁹ ~ 420 K for Si ions,²⁵ and ~ 490 K for Xe ions.²³ It is noted that both Weber's^{21, 23} and Matsunaga's¹⁶ work did not observe a significant difference in amorphization threshold between α and β SiC, whereas Inui reported ~ 50 K lower threshold temperature for α -SiC.¹⁷

Experimental

The materials chosen for this study were Cree Systems single crystal 6-H alpha SiC, Cercom hot pressed β -SiC, and Carborundum sintered α -SiC (Hexaloy). Four ion beams were selected for implantation in orientations both normal and parallel to the basal plane. Specifically, 3.6 MeV Fe and 0.56 MeV Si ions were implanted as single beams, with a dual beam of 1.8 MeV Cl and 1.0 MeV He. The silicon beam was produced in the NV-500 accelerator at the SMAC user facility at ORNL while the Fe, Cl and He ion beams were produced in the Triple Ion Beam Facility at ORNL. In both facilities, samples were heat sunk to a thermalizer block and temperatures measured using thermocouples either in contact with the thermalizer block or a backing to which the block was clamped. Typical ion beam fluxes were 7.1×10^{11} to 2.6×10^{12} ion/cm²-s, producing peak displacement damage rates of ~ 0.2 to 1×10^{-3} dpa/s. Cross sectional transmission electron microscopy was performed on all samples. The determination of amorphization was made by bright and dark field imaging and the presence of spot-free, diffuse diffraction patterns. The amorphized band was measured and compared with the TRIM-92²⁶ code to determine the critical dose for amorphization. The damage calculations were obtained assuming a sublattice-averaged threshold displacement energy of 40 eV. High resolution TEM was performed with a JEOL 4000 microscope, while conventional microscopy was performed using a Philips CM-12 microscope. Samples prepared for microscopy were mechanically thinned and ion milled with argon ions at 6 keV and 15° using a liquid nitrogen cooled stage during milling. A final mill was performed at 3 keV, 9° for 10 minutes at ambient temperature to reduce surface amorphization associated with argon implantation.

Two TEM discs each of Carborundum sintered α -SiC (Hexaloy) and hot pressed β -SiC produced by Cercom were irradiated in the target region of the High Flux Isotope Reactor in an aluminum capsule that was in direct contact with the reactor coolant water, producing a specimen irradiation temperature of ~ 340 K. The fast neutron ($E > 0.1$ MeV) fluence was $\sim 1.5 \times 10^{26}/\text{m}^2$ corresponding to a displacement level of ~ 15 dpa at a damage rate of $\sim 1 \times 10^{-6}$ dpa/s. Densities on these specimens were obtained using density gradient columns using mixtures of tetrabromomethane-methylene iodide or ethylene bromide-bromoform for the unirradiated and irradiated samples, respectively. The accuracy of the measurement was better than 0.001 g/cc and was found to be reproducible in repeated measurements and consistent for duplicate irradiated samples.

RESULTS AND DISCUSSION

Amorphization of the single crystal SiC due to 3.6 MeV Fe ion irradiation at room temperature is shown in Figure 1. The upper cross-sectional micrograph shows the highest fluence specimen (5.7×10^{15} ions/cm²), which amorphized over the entire range. The electron diffraction pattern for the amorphous region is inset into Fig. 1a. At a fluence of 2.0×10^{15} ions/cm² (Fig. 1b), amorphization only occurred over a ~ 0.7 μm wide band centered near the peak damage region. Amorphization was not observed at any depth at the lowest Fe ion fluence shown in Fig. 1c (1.4×10^{14} ions/cm²), which produced 0.05 dpa in

the peak damage region. Similar results were obtained for CVD SiC irradiated with Fe ions at the same three fluences.²⁵ No significant difference was found in the amorphization behavior of the Cree single crystal 6H-SiC, CVD SiC or Cercom hot pressed β -SiC. Figure 1 also demonstrates the swelling associated with the amorphization of SiC. By comparison of Fig. 1a (totally amorphous) with the non-amorphized sample (Fig. 1c), it is seen that amorphization has caused a significant volumetric expansion, with a resultant 0.27 μm movement of the free surface to the left. From this free surface "step-height" movement a volumetric swelling of $\sim 15\%$ was calculated due to the crystal-to-amorphous transformation.²⁵ Part of the original surface is missing from Fig. 1b so the volumetric expansion associated with the 0.7 μm amorphous band is not evident. From a comparison of the visible damage depths for the three Fe ion irradiation fluences, the amount of volumetric swelling associated with amorphization was calculated to be $\sim 17\%$.²⁵

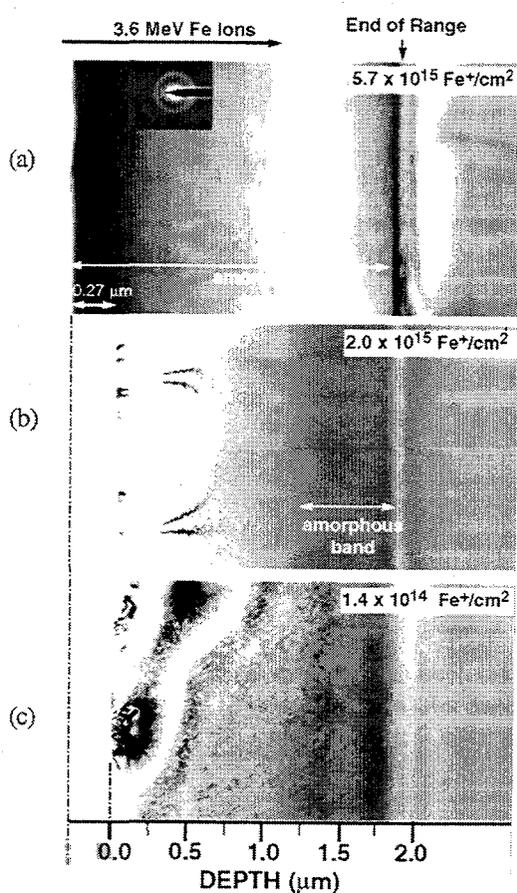


Figure 1. 3.6 MeV Fe ion amorphized SiC He ion single crystal 6H-SiC.

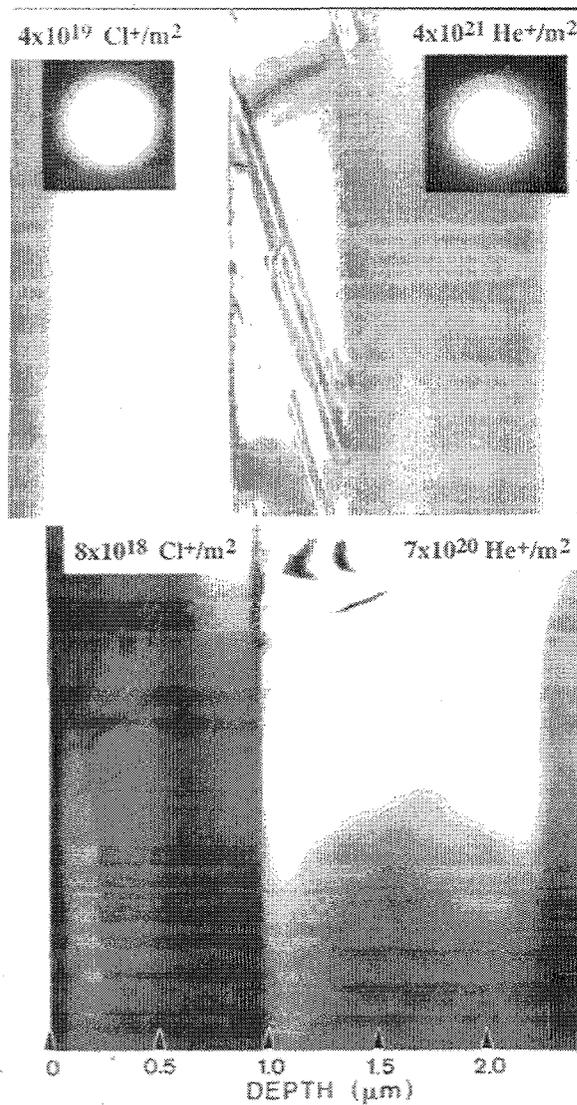


Figure 2. 1.8 MeV Cl and 1 MeV He ion amorphized β -SiC.

Figure 2 shows the effect of a room temperature simultaneous irradiation of 1.8 MeV Cl and 1.0 MeV He ions into a β -SiC sample.²⁵ The Cl beam had a damage range of nearly 1 μm , and the He beam produced damage up to ~ 2 μm . The upper micrograph of the figure is for the higher dose condition, and it is seen that amorphization has occurred over all but the near-surface regions of the Cl plus He ion irradiated region (~ 1.0 μm) and over the last third of the He ions irradiated region (maximum range ~ 2.4 μm due to density decrease associated with amorphization.) The lower ion doses shown in the bottom of this figure exhibit amorphization only near the end of the Cl and He ion ranges. Figure 2 also shows the effect of volumetric expansion, visible through the free surface movement. A measured swelling associated with amorphization of 17% ($\pm 3\%$) was measured for these specimens. The measured volume change associated with amorphization in SiC of 17% for the Fe, Cl and He ion irradiation is somewhat lower than the $22 \pm 3\%$ reported by Weber,⁵ but in good agreement with the value of 15% reported by McHargue.⁴

A matrix of samples was irradiated from room temperature to 400°C at several different Si ion fluences. Figure 3 shows examples of two different specimens that were irradiated to a Si ion fluence of 1.0×10^{15} ions/cm² at room temperature and 200°C. The room temperature specimen contained a band of material starting at about mid-range ($\sim 0.3 \mu\text{m}$) that appeared to be amorphous under ordinary bright field imaging conditions. However, through high resolution imaging a small volume fraction ($< 5\%$) of dispersed islands of crystalline material was observed at the beginning half of the "amorphous" band and was absent at the end of the ion range. These crystalline islands account for the weak diffraction spots seen in the diffuse rings of the upper diffraction pattern of Fig. 3 and indicate that the material was on the verge of total amorphization. The 200°C irradiated sample given in the lower part of Fig. 3 shows no evidence of amorphization for equal dose rates (1×10^{-3} dpa/s), indicating a temperature dependent dynamic recovery.

The measured threshold for amorphization as a function of temperature for the Si ions and the other ions is given in Figure 4. As discussed previously, the presence of Fe and Si ions did not have any measurable effect on the threshold amorphization dose, whereas Cl ions significantly reduced the threshold dose for amorphization in the implanted ion region.^{7, 25} This enhancement of amorphization is clearly demonstrated in Fig. 4 for the Cl irradiated material near room temperature. The remaining data in Fig. 4 were all taken from the unimplanted regions. Figure 4 shows a rapid increase in the critical amorphization dose with increasing temperature. At 420 K, no amorphization was observed in either the implanted or unimplanted regions at the highest Si ion dose studied, which was equivalent to 20 dpa.

These temperature dependent results are qualitatively consistent with previous studies which have used an in-situ technique for measuring the temperature dependent amorphization in SiC. For the in-situ approach, an irradiating beam is passed through, or partially through, a thin section of material with real-time observation of the microstructural evolution using TEM. Specifically, Kinoshita¹⁵ and Matsunaga¹⁶ used

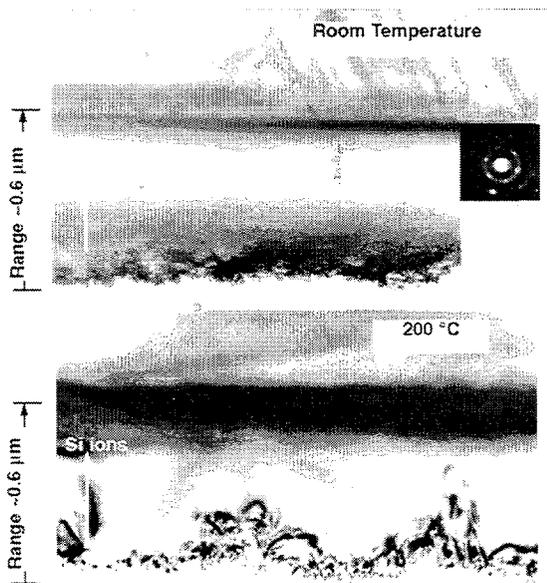


Figure 3. Cross sectional TEM of Si ion implanted SiC.

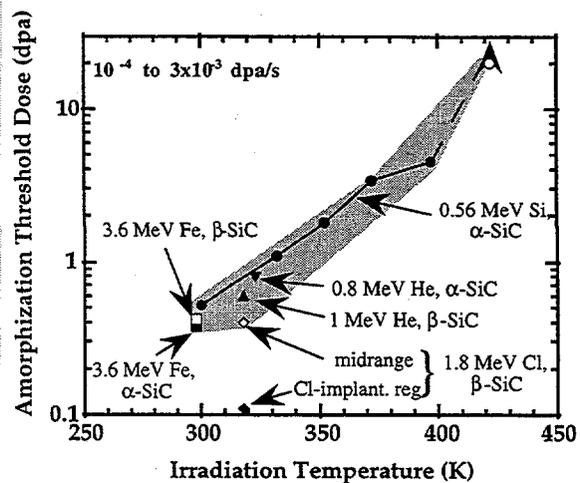


Figure 4. Amorphization threshold of SiC as a function of irradiation temperature.

1 MeV electrons while Inui¹⁷⁻¹⁹ has used 2 MeV electrons to both damage and image the samples. Weber^{5, 21-24} used a 1.5 MeV Xe beam in conjunction with TEM. In these examples Inui found the amorphization dose threshold below room temperature to be approximately 1 dpa for α -SiC¹⁷ and about 0.5 dpa for highly faulted β -SiC.¹⁸ Weber²¹ found the amorphization threshold near room temperature to be about 0.2 dpa for 1.5 MeV Xe ions. Figure 5 compares the data from the Si ion irradiations with the data from Inui and Weber, all of which were obtained at a damage rate near 1×10^{-3} dpa/s. All of the data in Fig. 5 have been plotted assuming an average displacement energy of 40 eV.

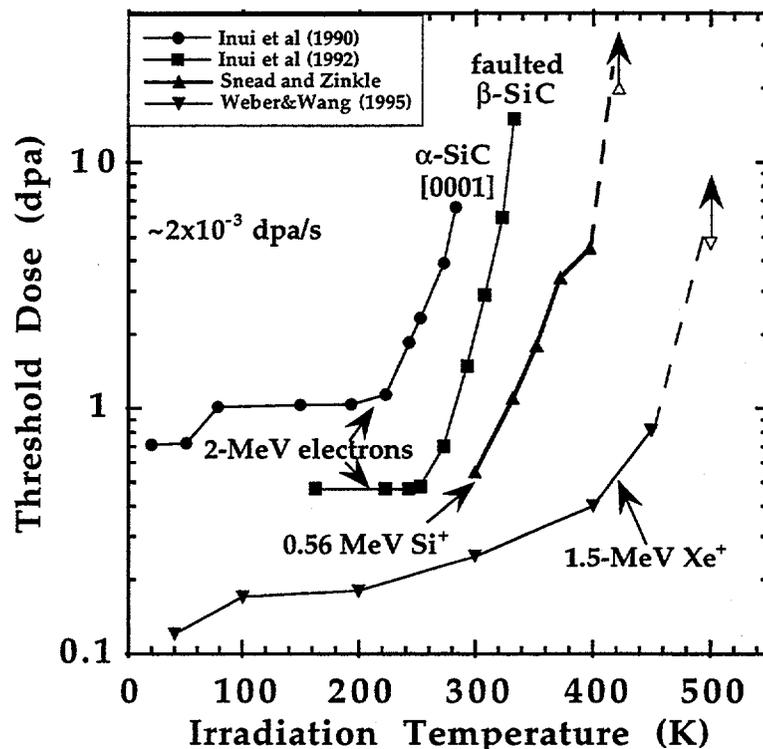


Figure 5. Temperature dependent critical dose for amorphization of SiC²⁵.

The results of Fig. 5 raise interesting questions regarding the mechanism for amorphization in SiC. Because amorphization has occurred under electron irradiation it is apparent that the SiC crystal, in the cubic or hexagonal form, can be driven amorphous by simple defect accumulation alone. The Xe ions, which can transfer much higher energies to the lattice atoms than the electrons (or the Si ions), are apparently capable of amorphizing SiC to a much higher temperature than electrons at the same damage rate. While the Xe ions will cause relatively large cascades compared to Si ions, no cascades can form for the low energy transfer electron irradiations. This supports the possibility of amorphization beginning within cascade structures which would also explain why the Xe ions produce amorphization to a higher temperature than the Si ions. Alternatively, it is possible that the heavier Xe ions simply create a higher concentration of "stable" interstitials than the Si ions or electrons on a dpa-normalized basis, thus driving the temperature threshold higher.

A potential problem in comparing the amorphization threshold data sets presented in Fig. 5 is whether the in-situ approach of Inui and Weber^{5, 17-19, 21-24} can be directly compared with the ex-situ measurements of the present study. In other words, are there technique-dependent effects occurring for one or both types of amorphization measurements. One example of an experimental artifact is demonstrated by observing a series of high resolution TEM images from the Si implanted, single crystal β -SiC. In general, there is a fairly rapid transition from crystalline material into the amorphous band (see Figs. 1-3). As expected from the damage gradient of the ion profile,⁷ this transition region is wider at the near-surface side of the

amorphization band compared to the end of the band (larger gradient in damage.) Figure 6 shows a HRTEM bright field image of the crystalline to amorphous transition region at the end of the amorphous band for a 6H-SiC sample irradiated with Si ions normal to the c-axis at 60°C. Fairly large areas of amorphous regions are seen, mixed into the crystal structure, which is consistent with the statistical nature of a point defect accumulation driven amorphization. Of particular interest in this micrograph is the shape and orientation of the amorphous areas, which appear to form as island structures with the major radii aligned parallel to the surface of the specimen. This phenomenon was observed in all amorphized specimens studied at both the beginning and end of the amorphized region from 60°C to 125°C for the Si ion irradiated specimens as well as the room temperature 6H SiC and CVD SiC specimens irradiated with 3.6 MeV Fe. These anisotropic amorphous structures occur when the beam is injected normal to the

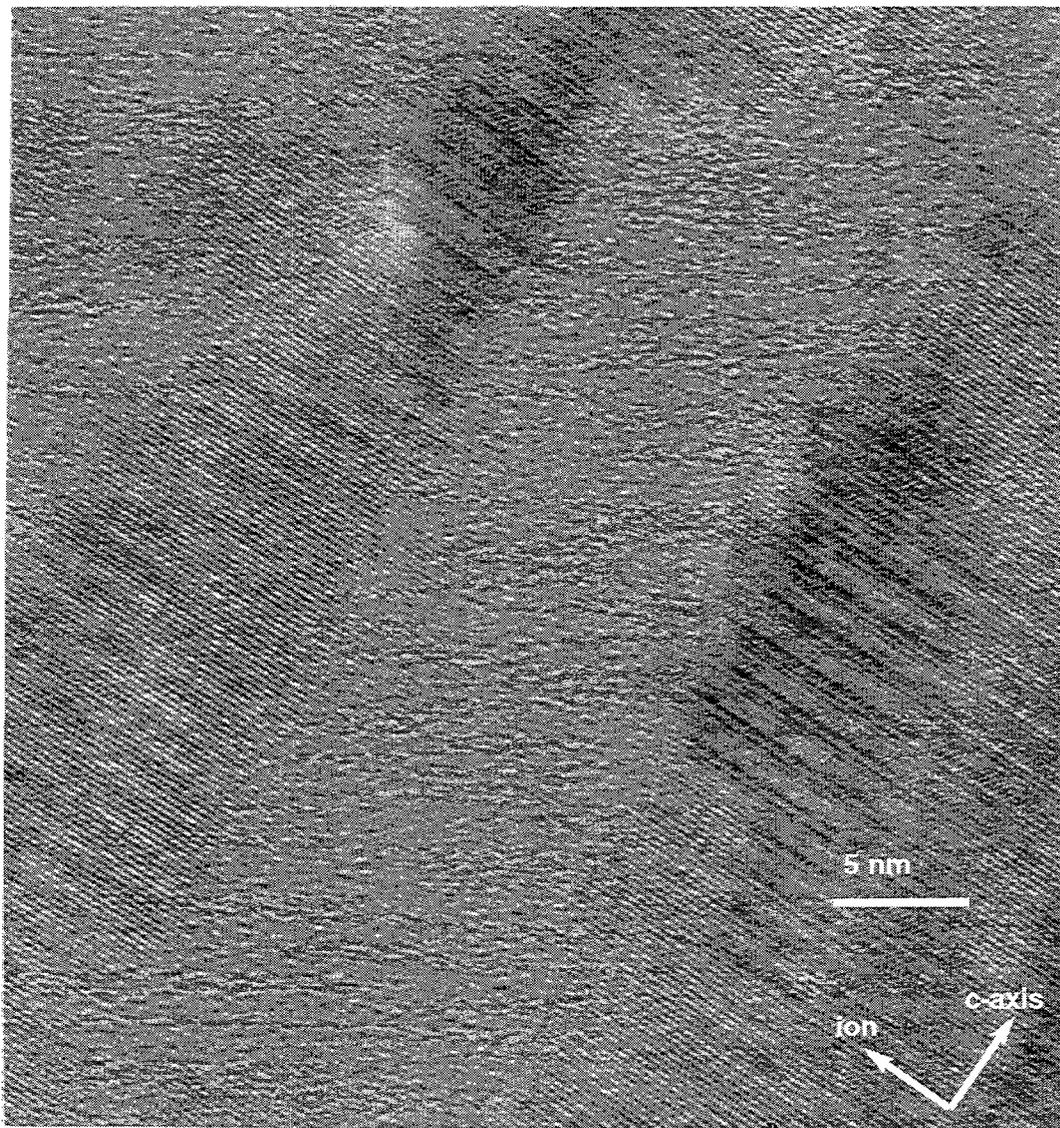


Figure 6. HRTEM image showing large island amorphous structures in 60°C Si Irradiated SiC.

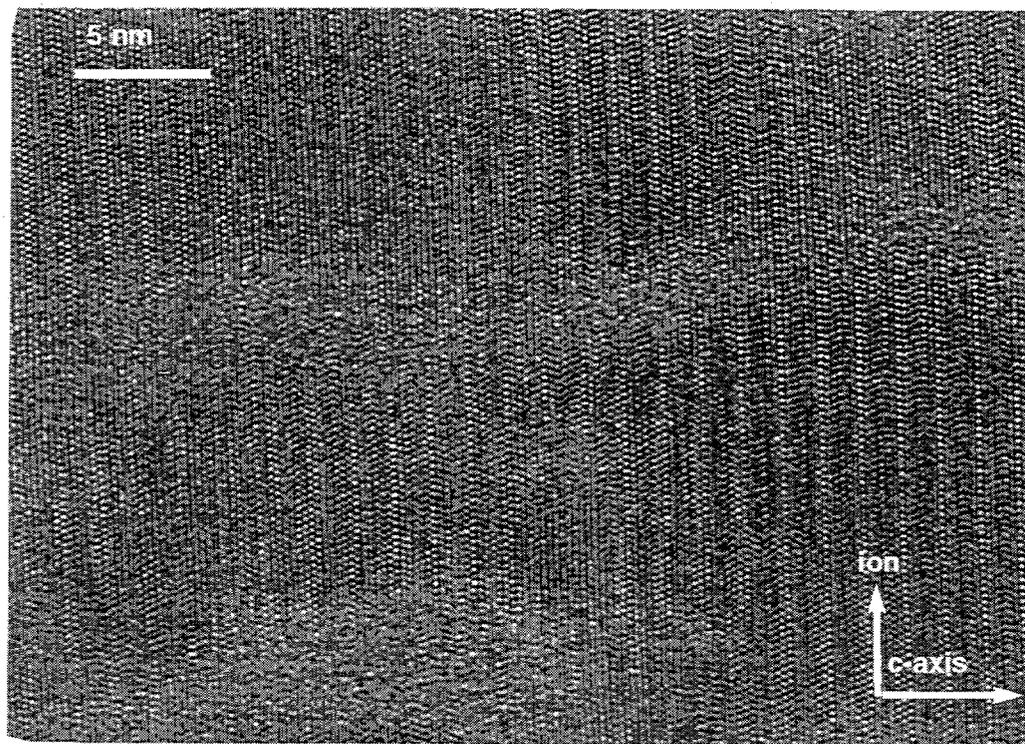


Figure 7. HRTEM image of Si irradiated SiC showing nucleation of island amorphous structures.

c-axis (as in Fig. 6) or parallel to the c-axis and therefore is not a crystallographic orientation phenomenon. Figure 7 shows another bright field micrograph illustrating that the formation of these islands occurs on a very fine scale, eventually growing and linking with other islands.

In order to quantify this effect, a series of micrographs were taken at the beginning and end of the amorphized band for specimens irradiated with Si ions at 60 and 125 °C. Amorphous islands were identified using the complete loss of lattice fringes as the criteria, and the major and minor radius of each amorphous pocket was measured. Typically, about 50 islands were measured from at least three separate areas on each specimen. Results of this analysis showed that the aspect ratio (major/minor radius) was fairly constant independent of the temperature and location of the amorphous band. Specifically, the aspect ratio was found to be 3.3 and 3.1 for the beginning and end of the amorphous band for the 60°C Si ion irradiated material and 3.2 and 3.6 at the beginning and end of the amorphous band for the 125°C irradiated material, respectively.

The anisotropic geometry of these islands relative to the surface of the specimens raises the question of whether there is a strain effect on the evolution of the amorphous structure. There is obviously a large

lattice strain induced in the crystal to amorphous transition both from point defect accumulation and the amorphous pockets. Also, the presence of the free surface should affect the nature of the strain field. The obvious question from these results is whether the formation and connection of the amorphous pockets are being affected in a substantial way by this phenomenon. For example, it has been shown that the amorphization threshold is reduced in Cl ion implanted SiC.⁷ Amorphous regions could nucleate first in such an implanted region causing a large stress field due to the associated volumetric expansion and contribute to drive the surrounding material amorphous. The presence of the free surface and/or the expansion of the material in a direction towards the surface could affect this strain field and the subsequent amorphization.

Chemical rate theory models have been developed by several authors which can be fitted to experimental amorphization data. However, these models contain numerous parameters which are not well known and generally do not have analytical solutions. More simplistic analytical models have been developed.²¹ However, these models have yielded unrealistically low activation energies for amorphization (below 0.1 eV), considering that the recovery processes occur near or above room temperature.

The data of Fig. 5 have been analyzed with an analytical approximation based on the rate theory model of Motta and Olander²⁷ developed for electron induced amorphization of intermetallic precipitates in zircaloy. This model is analogous to well known rate theory results for point defect concentrations in irradiated metals and is based on the free energy increase associated with point defect accumulation, and is therefore not appropriate if in-cascade amorphization occurs. The assumption of point defect driven amorphization is supported by the molecular dynamics simulations²⁸ and high-resolution TEM observations^{18,23} which indicate that amorphization in SiC proceeds gradually by point defect accumulation processes, as opposed to heterogeneous in-cascade amorphization. There is no evidence for the spontaneous creation of amorphous zones in displacement cascades in irradiated SiC,^{23,28} though this should still be considered an open question.

Specifically, the model²⁷ as applied to amorphization in SiC uses the temperature dependent term for the critical damage for amorphization (see Fig. 5) which is given by (1) :

$$(1) \quad d(T) = \frac{A}{\sqrt{\phi}} e^{-\frac{E_m^i}{2kT}} \quad (2) \quad dpa_{crit} = d(T) + a$$

The critical dose for amorphization as given in Fig. 5 can be described by the sum of this temperature dependent term and the temperature independent term obtained at low temperatures as shown in equation (2).

The interstitial migration energy, E_i^m , is found from the slope of the experimental amorphization data, as plotted in Figure 8. The interstitial migration energy for the silicon ion irradiation was calculated to be 0.56 eV and is in agreement with the data sets of Inui¹⁷ and Weber⁵ on α -SiC. The data by Inui¹⁹ on highly faulted β -SiC is significantly higher than this value, however. All the data sets of this figure are at approximately the same displacement rate ($\sim 2 \times 10^{-3}$ dpa/s) rendering the pre-exponential term in equation (1) constant.

The physical explanation for the temperature-dependent behavior of the Motta and Olander amorphization model is associated with interstitial migration that assists restoration of chemical order in the lattice. We assume that amorphization occurs via chemical disordering on both the Si and C sublattices in SiC. Experimental measurement of the threshold energy to induce amorphization in SiC during cryogenic temperature electron irradiation is needed to determine the validity of this assumption. The migration energy of 0.56 eV is tentatively attributed to the slower interstitial specie (presumably Si). As shown in Fig. 5, both α -SiC electron¹⁷ and β -SiC Xe ion²¹ irradiation studies observed a step in the amorphization dose at low temperature (~ 70 K) which may be due to migration of the faster interstitial specie, or perhaps recovery of close pair Frenkel defects.

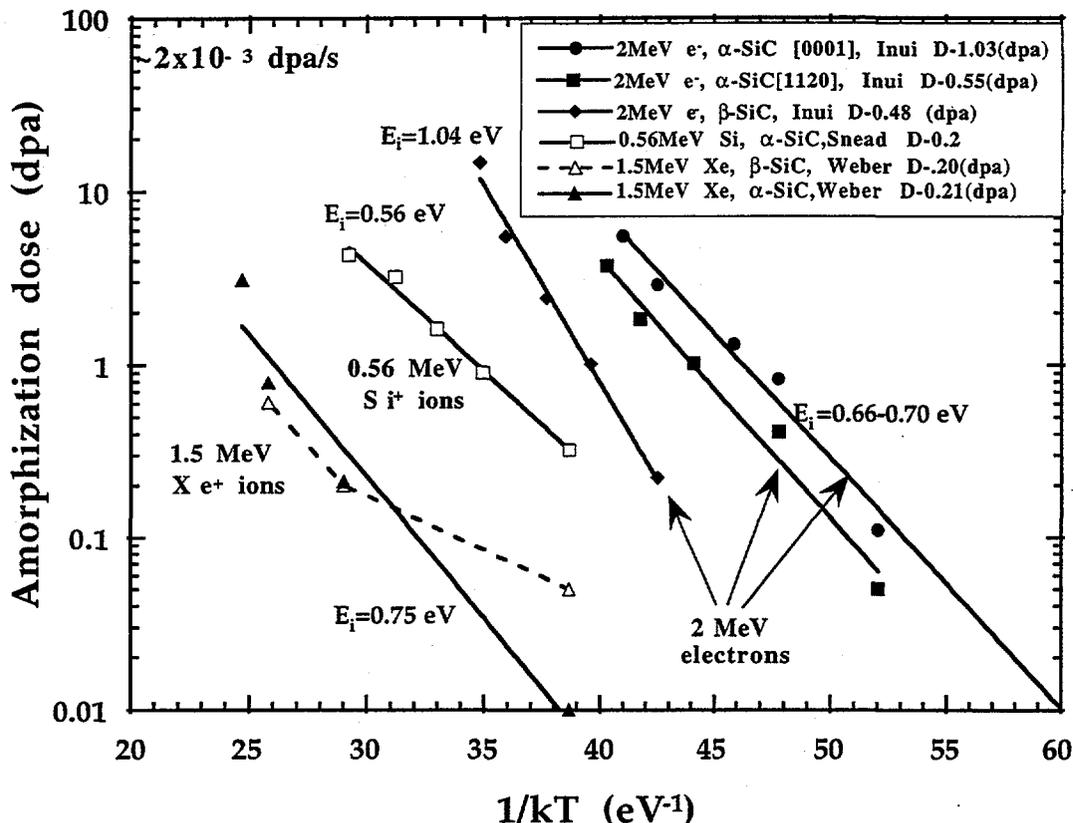


Figure 8. Interstitial migration energy determination for amorphization of SiC.

Referring again to Fig. 5, it is seen that for all the data sets used, amorphization has occurred in the temperature independent region at a displacement dose (d_0) of less than 1 dpa. The temperature above which amorphization is not observed is ~ 300 K for electron irradiated material, ~ 420 K for Si ions, and ~ 500 K for Xe ions. The SiC material which was neutron irradiated in this study to a dose of approximately 15 dpa yielded density changes of 10.1% for the Cercom β -SiC and 11.1% for the Carborundum α -SiC. While both values are significantly higher than expected for simple "point defect" swelling from the previous neutron irradiation data for swelling in silicon carbide, it is less than the 15% swelling level reported for complete amorphization. However, given this large swelling which has occurred it is likely that the samples include pockets of amorphization. The neutron irradiation temperature for these samples was ~ 350 K.

Given that the fast neutron irradiation of SiC should produce a damage spectrum closer to the ion irradiated material than the simple point defect only spectrum of the electron irradiation it might be expected that complete amorphization should have occurred during neutron irradiation at the ~ 350 K neutron irradiation. The apparent partial retention of the crystal structure for the neutron irradiated SiC specimens can be explained with the above rate theory model by noting that the damage rate for the neutron irradiation was $\sim 1 \times 10^{-6}$ dpa/s which is three orders of magnitude lower than the Si ion irradiation. Using the 0.56 eV migration energy calculated from the Si ion data and using the respective damage rates, equation (2) can be used to plot the calculated threshold dose profile for amorphization at two different damage rates, as shown in Figure 9. This plot includes the silicon ion irradiation data along with its calculated profile and the data point for which amorphization was not observed (20 dpa and 420 K.)

The threshold dose profile for the fission neutron damage rate of $\sim 1 \times 10^{-6}$ dpa/s is calculated with the derived 0.56 eV interstitial migration energy, and migration energies ± 0.05 eV from this value, to

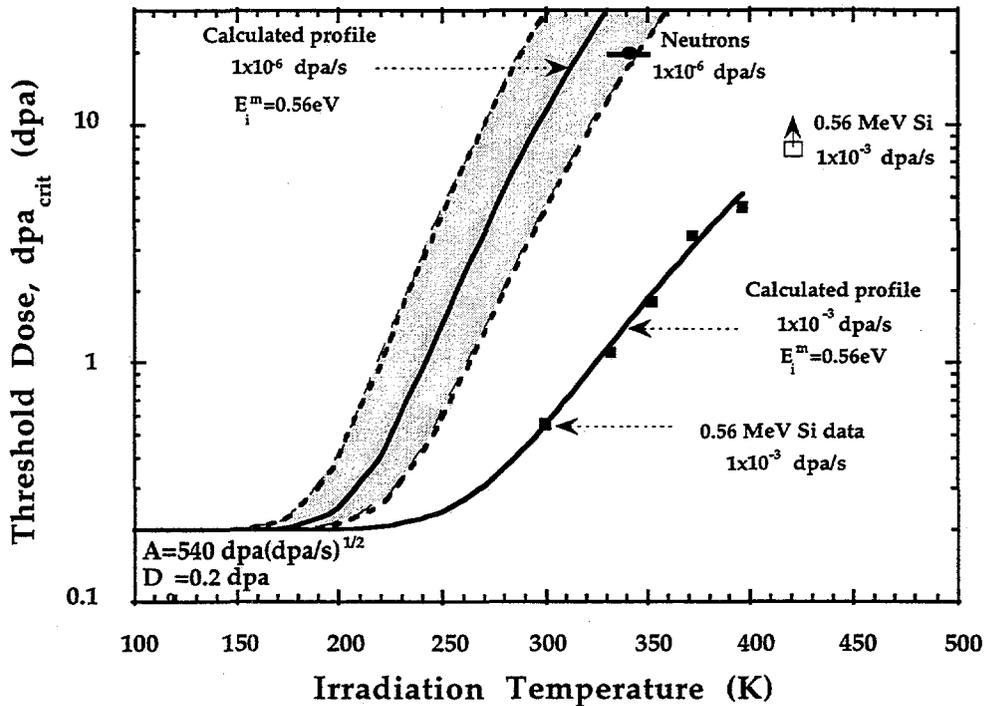


Figure 9. Predicted flux dependence of the critical dose for amorphization of SiC.

demonstrate sensitivity to this variable. Given the nature of the calculation of the interstitial migration energy, such an uncertainty is conceivable. However, assuming a migration energy of 0.56 eV, it is seen that for the ~ 340 K fission neutron irradiation temperature and $\sim 1 \times 10^{-6}$ dpa/s damage rate, it would require between 30 and 50 dpa to completely amorphize the SiC. For comparison it would require < 2 dpa at the Si ion dose rate of $\sim 1 \times 10^{-3}$ dpa/s. While this model appears to adequately describe the temperature dependence of amorphization at moderate temperatures, it is not entirely satisfying in that it does not exhibit the "critical temperature" asymptotic behavior which is suggested by the data. As the temperature is increased and the exponential term approaches unity, the critical dose for amorphization approaches a constant value of $d_o + A / \sqrt{\phi}$. Modification of the Motta and Olander model is necessary to correct this nonphysical predicted amorphization behavior at very high temperatures. The inadequacy of the present model at very high temperatures is seen by noting the unamorphized Si ion irradiated data point irradiated to 20 dpa at 420 K. The extrapolated Si ion model would fall to the right of this data point (implying the material should have amorphized) and continue to flatten towards a constant value just mentioned.

CONCLUSIONS

- (1) There is apparently no PKA spectrum effect in this study on the low temperature (temperature independent) amorphization of SiC when irradiating with He, Si, Cl and Fe atoms. However, published low temperature 2 MeV electron and 1.5 MeV Xe irradiations suggest that the low-temperature amorphization dose decreases by about a factor of 5 as the irradiation source is changed from electrons to ions.

- (2) The critical temperature for amorphization of SiC with 0.56 MeV Si of ~420 K shows an exponential dependence with an apparent critical temperature threshold above which amorphization dose not take place at a damage rate of $\sim 2 \times 10^{-3}$ dpa/s. This compares with the published amorphization temperature thresholds for 2 MeV electron and 1 MeV Xe ions of ~300 K and 500 K respectively at a similar damage rate.
- (3) A possible experimental artifact was observed which may affect the calculation of critical dose for amorphization in ion implanted specimens. Specifically, amorphized islands were seen to form with their major radii parallel to the free surface of the specimen. These islands have an aspect ratio of ~3 which are independent of temperature and crystal orientation. Formation of these islands raises the question of whether surface or strain effects alter the apparent amorphization thresholds in the bulk irradiation method presented here and the thin foil in-situ method which has been used by other workers. A determination of the adequacy of this implantation technique for determining amorphization threshold should be conducted and compared with the in-situ "thin foil" technique.
- (4) A swelling of approximately 10% was measured for sintered α -SiC and hot pressed β -SiC neutron irradiated to 15 dpa at ~350 K. This amount of swelling is about a factor of three higher than previously reported for neutron irradiation at similar temperatures. However, this amount of swelling is less than the ~15 to 20% level of swelling observed in completely amorphized ion-irradiated material. It is likely that partial amorphization characterized by amorphous islands is present in the material. XRDA and TEM analysis will be performed to determine the extent of amorphization in the room temperature irradiated samples.
- (5) A model for the critical dose for amorphization based on the free energy increase associated with point defect accumulation in metals is applied to the temperature dependent critical dose for amorphization. An interstitial migration energy of 0.56 eV was calculated based on the Si ion (10^{-3} dpa/s) irradiations and adequately describes the transition from the temperature independent to temperature dependent regimes of amorphization. However, the model does not adequately describe the higher temperature asymptotic amorphization dose behavior. Extrapolation of this model to the dose rate for the neutron irradiated samples (10^{-6} dpa/s) successfully predicts the observed lack of complete amorphization in these specimens. Assuming this dose rate and the interstitial migration energy of 0.56 eV, it would require between 30 and 50 dpa to have reached total amorphization for the material neutron irradiated at ~340 K. Presently, a neutron irradiation capsule identical to the one used in this study is in the HFIR reactor and will achieve 60 dpa. Examination of the SiC specimens from this capsule should provide interesting information on the validity of the fitted rate-controlling interstitial migration energy, and whether SiC can be amorphized during high-dose neutron irradiation at ~340 K.

REFERENCES

1. W. Primak, L.H. Fuchs, and P.P. Day, *Phys. Rev.*, **103**(5), 1184-1192 (1956).
2. R. Blackstone, *J. Nucl. Mater.*, **39**, 319-322 (1971).
3. C.W. White et al., *Nucl. Instrum. Meth. B*, **80/81**, 41-146 (1989).
4. C.J. McHargue and J.M. Williams, *Nucl. Instrum. Meth.*, **B80/81**, 889-894 (1993).
5. W.J. Weber et al., accepted for publication, *J. Nucl. Mater.* (1996).
6. L.L. Snead, S.J. Zinkle, and D. Steiner, *J. Nucl. Mater.*, **191-194**, 560-565 (1992).
7. L.L. Snead and S.J. Zinkle, *MRS Symposium #373 on Microstructure of Irradiated Materials* (1995) Materials Research Society, Pittsburgh, PA., I. M. Robertson et al., eds.
8. R.R. Hart, H.L. Dunlap, and O.J. March, *Rad. Effects*, **9**, 261-266 (1971).
9. J.M. Williams, C.J. McHargue, and B.R. Appleton, *Nucl. Instrum. Meth.*, **209/210**, 317-323 (1983).
10. J. Spitznagel et al., *Nucl. Instrum. Meth.*, **B16**, 237-243 (1986).
11. J.A. Edmond et al., *J. Mater. Res.*, **3**, 321-328 (1988).
12. N.G. Chechenin et al., *Nucl. Instrum. Meth.*, **B65**, 341-344 (1992).
13. L.L. Horton et al., *Nucl. Instrum. Meth.*, **B65**, 345-351 (1992).
14. R.G. Vardiman, *Mat. Res. Soc. Symp.*, 1989.

15. C. Kinoshita et al., in *Proceedings of the 11th International Conference on Electron Microscopy* (1986), Japanese Society of Electron Microscopy.
16. A. Matsunaga et al., *J. Nucl. Mater.*, **179-181**, 457-460 (1991).
17. H. Inui, H. Mori, and H. Fujita, *Phil. Mag.*, **B 61**, 107 (1990).
18. H. Inui, H. Mori, and T. Sakata, *Phil. Mag.*, **B 66(6)**, 737-748 (1992).
19. H. Inui et al., *Phil. Mag.*, **B 65**, 1 (1992).
20. S.D. Harrison and J.C. Corelli, *J. Nucl. Mater.*, **99**, 203-212 (1981).
21. W.J. Weber and L.M. Wang, *Nucl. Instrum. Meth.*, **B 106**, 298-302 (1995).
22. W.J. Weber and N. Yu, *Accepted Nucl. Instrum. Methods B*, (1996).
23. W.J. Weber, L.M. Wang, and N. Yu, *Nucl. Instrum. Meth.*, **B 116**, 322-326 (1996).
24. W.J. Weber. *13th International Cong. on Defects in Insulating Materials* (1996).
25. S.J. Zinkle and L.L. Snead, *Nucl. Instrum. Meth. B*, **116**, 92-101 (1996).
26. J.F. Ziegler, J.P. Biersak, and U.L. Littmark, Pergamon, New York (1985).
27. A.T. Motta and D.R. Olander, *Acta Metta. Mater.*, **38**, 2175-2185 (1990).
28. T.D. de la Rubia, M.-J. Caturla, and M. Tobin. in *Microstructure of Irradiated Materials* (1995) Boston, MA., I. M. Robertson et al., eds.